

Multiple Frequency Surface Resistance Measurement Technique Using a Multimode TE_{01n} Cylindrical Cavity on a TlBaCaCuO Superconducting Film

J. S. Deroy, A. J. Drehman, M. J. Suscavage, R. J. Andrews, E. Cohen, N. Ianno, and D. Thompson

Abstract—A multiple frequency surface resistance measurement technique for high-temperature superconductors has been developed using a modified multimode cylindrical cavity made of niobium operating in the superconducting state. The surface resistance vs. frequency relationship of TlBaCaCuO high-temperature superconducting thin films was determined over the frequency range of 50–100 GHz with an extrapolation down to 1 GHz. The TE_{01n} ($n = 1, 2,$ and 3) modes of the cylindrical niobium waveguide cavity had resonances at 52, 71, and 95 GHz and the measured surface resistances at these frequencies were 5.1, 8.5, and 21.1 milli-ohms, respectively.

I. INTRODUCTION

FREQUENCY dependent surface resistances of high-temperature superconducting materials provide fundamental information for microwave and millimeter wave circuit design. Our purpose was to develop a technique to measure the surface resistance of high-temperature superconductors in the 50–100 GHz range using a single high- Q TE_{01n} cylindrical cavity for more than one frequency. Stripline resonators are typically used to make frequency dependent surface resistance measurements and work well below 50 GHz [1], but this becomes more difficult above frequencies at which coaxial connectors are no longer available. TE_{01n} mode cavity measurements have yielded surface resistance results at frequencies greater than 50 GHz, but these were single frequency measurements and required separate cavities for each frequency [2], [3]. In this letter we demonstrate the use of a single TE_{01n} mode cylindrical waveguide reflection cavity to make frequency dependent surface resistance measurements at 52, 71, and 95 GHz. This allowed us to take all three frequency data points during one temperature cycling of the system. A superconducting niobium TE_{01n} mode cavity ($n = 1, 2, 3$ and with resonances at 52, 71, and 95 GHz) was used to determine the frequency dependent surface resistance of a laser deposited Thallium–Barium–Calcium–Copper Oxide (TlBaCaCuO) thin film. This method can be extended to other high-temperature superconductors.

Manuscript received July 22, 1992. This work was supported by the Rome Laboratory, Laboratory Director's Fund.

J. S. Deroy, A. J. Drehman, M. J. Suscavage, and R. J. Andrews are with the Rome Laboratory, Electromagnetics and Reliability Directorate, Hanscom AFB, MA 01731-5000.

E. Cohen is with ARCON, Inc., 260 Bearhill Road, Waltham, MA 02154.

N. Ianno is with the Department of Electrical Engineering, University of Nebraska, Lincoln, NE 68508.

D. Thompson is with SEI, Inglewood, CA.

IEEE Log Number 9204281.

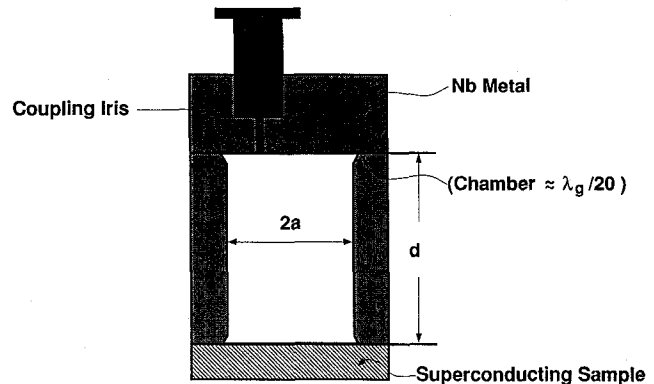


Fig. 1. Cavity cross-section and placement of the superconducting sample during measurement. The cavity was a reflection type cavity and its dimensions where $a = 4.213$ mm and $d = 5.31$ mm.

II. CAVITY MODES

A frequency dependent surface resistance measurement can be made using the TE_{01n} modes ($n = 1, 2,$ and 3) in a cylindrical waveguide cavity. The TE_{01n} modes were chosen because of their unique field symmetry where no surface current flows between the cylindrical side walls and the end walls of the cavity [4]. Therefore, by replacing one end wall of the cavity with the superconducting film, no error is introduced in the loss measurement through discontinuities between the end wall and the cylindrical side wall. Since the TE_{01n} mode can be degenerate with the TM_{11n} mode, a chamfer was cut in the side wall of the cavity as shown in Fig 1. The chamfer affects the TM field distribution, while having a negligible effect on the TE field distribution. This removes the degeneracy. Another method for removing the degeneracy has been reported by B. Mayer *et al.* [5]. For this work, the cavity was built from niobium (Nb) because this material is known to have very low loss in the superconducting state. All the existing TE and TM modes of the cavity were determined [6] and a cylindrical cavity mode chart of the TE_{011} , TE_{012} , and TE_{013} modes and their nearest neighboring modes was constructed. This is shown in Fig. 2. In our design, we chose $(2a/d)^2$ to be 2.4 with $a = 4.213$ mm and coupling iris of about 1 mm in diameter. As seen from this figure, the neighbors of the TE_{011} and the TE_{012} modes are far apart, but the TE_{013} is only 400 MHz from the unwanted TE_{222} mode. Since we are using a niobium cavity whose quality factor was measured to be 8×10^5 at 95 GHz with a bandwidth on the

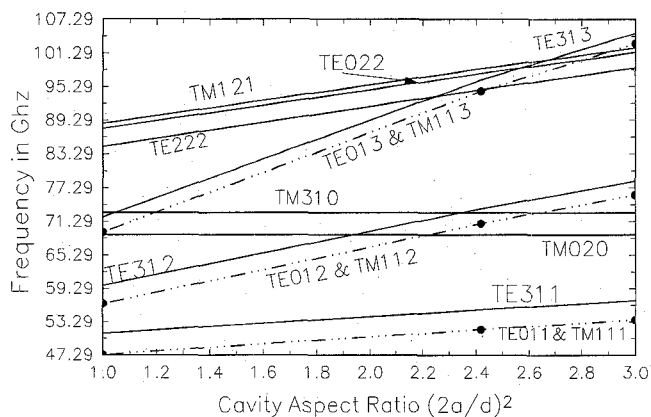


Fig. 2. Mode chart for the cylindrical cavity used to make the surface resistance measurements.

order of 120 KHz, the two modes can easily be resolved.

The surface resistance of the high-temperature superconducting end wall in the Nb cavity can be determined from a measurement of the cavity Q . The relationship between the Q of the cavity and the absolute value of the surface resistance R_{TI} for the thallium film is given by [7]

$$R_{TI} = \mu\pi f_o \left(\frac{1}{Q_{TI}} - \frac{1}{Q_{Nb}} \right) \frac{\int_{vol} |H|^2 \rho d\rho d\phi dz}{\int_{end} |H \times z_o|^2 \rho d\rho d\phi} + R_{Nb}, \quad (1)$$

where ρ, ϕ, z , denote the axes of the cylindrical coordinate system with their respective unit vectors ρ_o, ϕ_o, z_o . In (1), f_o is the resonant frequency, H is the magnetic field for the $(0,1,n)$ mode, Q_{TI} and Q_{Nb} are the values of Q for the cavity with the thallium film as an end wall and with a Nb film as an end wall respectively, and R_{Nb} is the surface resistance of the niobium.

III. MEASUREMENT AND RESULTS

The cavity Q was determined by a standard measurement which is based on measuring the resonant frequency and bandwidth of the cavity at the 3 db points [4], [8]. WR-15 waveguide, which has a cut-off frequency of 39.86 GHz, was used from the cavity to the top of the cryostat and a WR-10 to WR-15 waveguide transition was used between the source and the cryostat to make the 95-GHz measurement. The cavity Q was measured with the thallium film as the end wall and with a niobium film as the end wall. The surface resistance of the thallium film was then calculated from (1). The calculated surface resistance for the thallium film was 5.2, 8.5, and 21.2 milli-ohms at 52, 71, and 95 GHz, respectively. Since we were using a niobium cavity, the measurements were performed at 4.2 K. These surface resistance values are at least a factor of 15 greater than the minimum resolution of our system. The $Tl_2Ba_2CaCu_2O_x$ film on $LaAlO_3$ substrate had a T_c of 105 K. The film was 8000-Å thick, which is approximately four times the estimated London penetration depth. Therefore, the losses due to the finite thickness of the film were considered to be negligible. The results of the measurement are plotted in Fig. 3. The dashed line in the figure is a least-squares best fit to the data assuming a f^2 dependence for the surface resistance. The f^2 dependence of the loss is

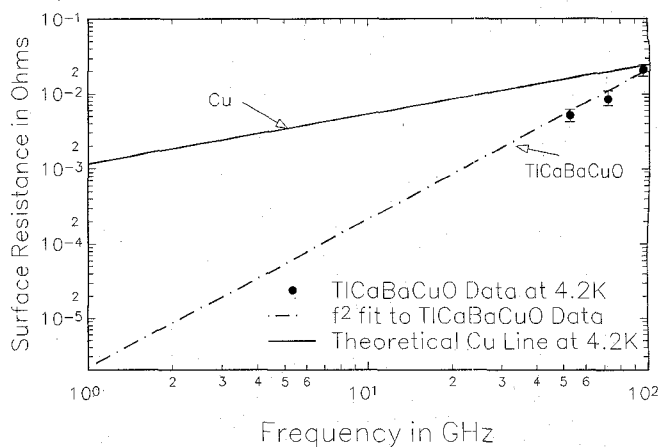


Fig. 3. Surface resistance versus frequency relationship for the measured TlBaCaCuO superconducting film.

predicted by the two-fluid model for the surface impedance of a superconductor [9]. The curve is extrapolated down to 1 GHz to estimate the superconducting material's performance at lower frequencies. Note that the surface resistance of the superconducting material is approximately equal to that of copper at 100 GHz.

IV. CONCLUSION

The surface resistance of a superconducting thallium film was measured using the TE_{01n} ($n = 1, 2, \text{ and } 3$) cylindrical cavity technique. To date, this is the simplest most reliable technique for measuring surface resistance of superconductors in the frequency range of 50 to 100 GHz. The agreement between the data and the f^2 dependence of the two-fluid model gives further credibility to this technique.

ACKNOWLEDGMENT

The authors wish to thank Dr. B. Tomasic, Dr. S. Mittleman, Dr. P. Carr, and Dr. R. Mailloux for their comments on this manuscript and Dr. G. Dresselhaus for his encouragement.

REFERENCES

- [1] D. E. Oates, A. C. Anderson, C. C. Chin, J. S. Dero, G. Dresselhaus, and M. Dresselhaus, "Surface-impedance measurement of superconducting NBN films," *Phys. Rev. Ser. B*, vol. 43, no. 10, pp. 7655-7663, Apr. 1991.
- [2] T. W. Button and N. McN. Alford, "High- Q $YBa_2Cu_3O_x$ cavities," *Appl. Phys. Lett.*, vol. 60, no. 11, Mar. 16, 1992.
- [3] L. Drabek, K. Holczer, G. Gruner, and D. J. Scalapino, "Ohmic and radiation losses in superconducting films," *J. Appl. Phys.*, vol. 68, no. 2, July 15, 1990.
- [4] E. L. Ginzton, *Microwave Measurements*. New York: McGraw-Hill, 1957, ch. 9, Appendix A.
- [5] B. Mayer, A. Reccius, and R. Knelch, "Conical cavity for surface resistance measurements of high-temperature superconductors," *IEEE Trans. Microwave Theory Tech.*, vol. 40, pp. 228-235, Feb. 1992.
- [6] R. E. Collins, *Foundations for Microwave For Engineering*. New York: McGraw-Hill, 1966, ch. 3.
- [7] J. D. Jackson, *Classical Electrodynamics*, 2nd ed. New York: Wiley, 1975, ch. 8.
- [8] M. Shucher and J. Fox, *Handbook of Microwave Measurements*. New York: Polytechnic, 1963, vol. 2, ch. 8.
- [9] T. Van Duzer and C. W. Turner, *Principles of Superconducting Devices and Circuits*. New York: Elsevier, 1981, pp. 128-130.